



MadDM

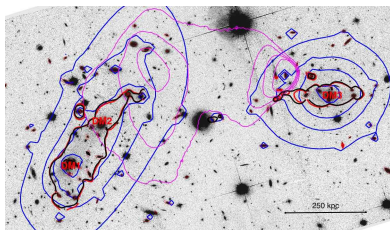
A computational tool of DM relic abundance and direct detection

M. Backovic, F. Maltoni, A. Martini, Mat McCaskey, K.C. Kong, G. Mohlabeng

The mystery of dark matter

Observations

- galaxies rotational curves.
- Bullet cluster.
- Cosmic microwave background :
 $\Omega_{DM} h^2 = 0.1199 \pm 0.0027.$



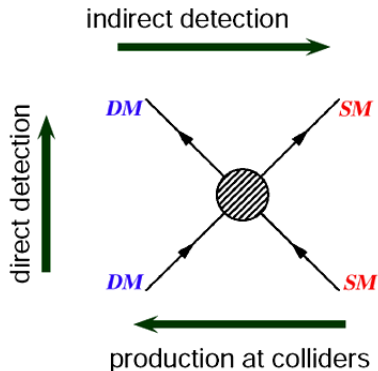
Candidates

- Heavy neutrinos
- Axions (Peccei-Quinn): CP violation in QCD.
- LSP (SUSY) : neutralino $\tilde{\chi}_1^0$, gravitino (supergravity)
- LKP (Kaluza-Klein) : extra-dimensions theories.
- ...

DM detection opportunities

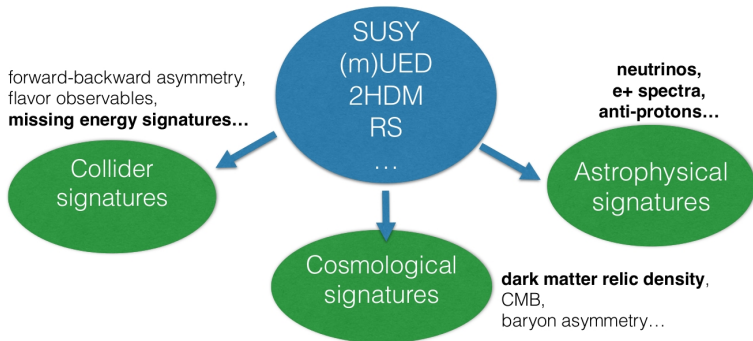
Three possibilities

- Indirect detection: FERMI-LAT, AMS-02.
- Direct detection: XENON100, LUX, CDMS.
- Production at collider: LHC.



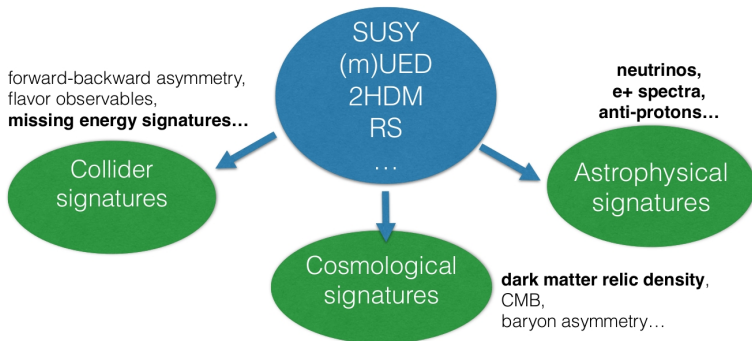
BSM Physics in the LHC era

- If new physics exists, **we don't have a good sense of the scale yet!**
- The result is a vast number of possibilities and many approaches to measure them.



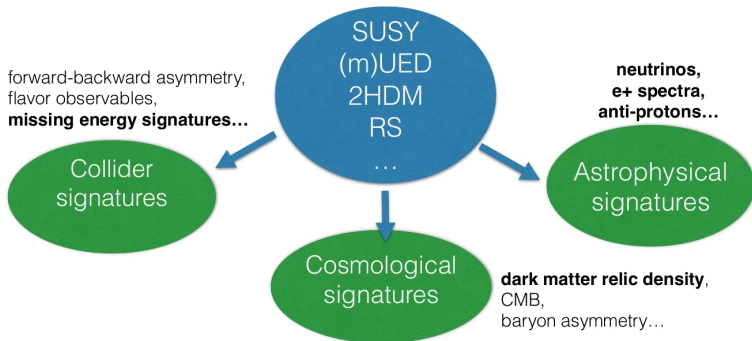
BSM Physics in the LHC era

- We have good “hints” that there is BSM physics out there - **dark matter is a good example!**
- Important to look for DM both at colliders and in the galaxy!

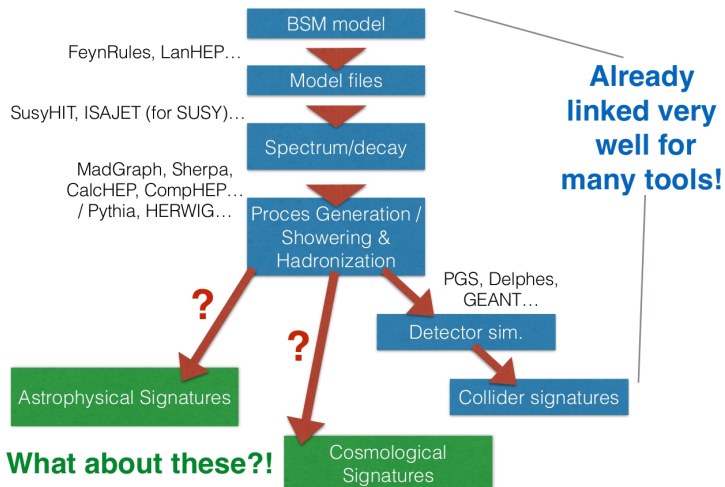


BSM Physics in the LHC era

- Models of new physics grew complex - Many fields, many parameters, many signatures.
- Much relies on the numerical tools nowadays.

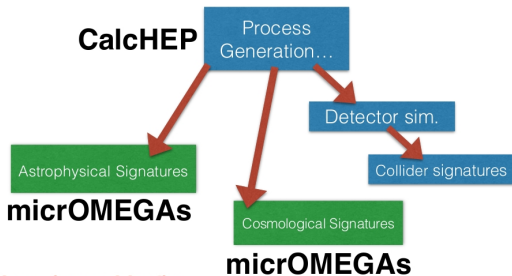


BSM Tools in the LHC era



DM Tools in the LHC era

Currently **CalcHEP/micrOMEGAs** (Comput.Phys.Commun. 149 (2002) 103-120) is the only tool available which can calculate Collider + Astrophysical + Cosmological signatures in a **generic model**.



Some limitations (not critical):

- 4 point interaction in SU(3) needs a special treatment. FeynRules does not generate correct model files for CalcHEP/micrOMEGAs.
- Practically limited by the fact that CalcHEP takes a long time to calculate collider processes beyond 4-6 final state particles.

DM tools in the LHC era

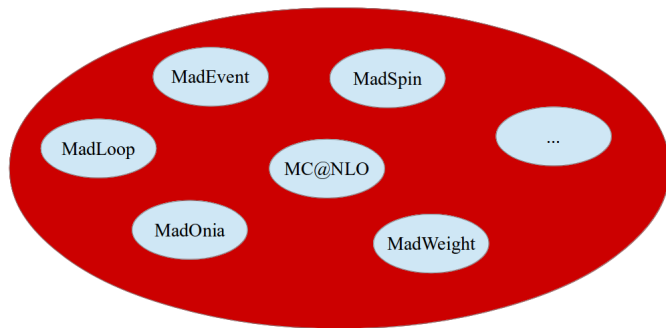
Tools

- **darkSUSY** is a popular, although model specific DM tool (JCAP 0407 (2004) 008).
- SUSY contains a lot of generic DM features like **co-annihilations and resonant annihilations** => darkSUSY has a lot of useful technology for DM phenomenology in a generic model.
- darkSUSY can be 'hacked' to include other models, but this is not trivial !
- Many other (model specific) tools exist: Isatools, SSARD, Drees, Roszkowski...

MADGRAPH5

Since 1994, MADGRAPH has grown into a powerful collider phenomenology framework...

MADGRAPH 5



... but no capability to calculate astrophysical and cosmological signatures in models which contain dark matter candidates.

MadDM

- First step to extend MadGraph5 capabilities to dark matter phenomenology
- Built on top of existing MadGraph5 architecture - inherits all of the existing MadGraph structure.

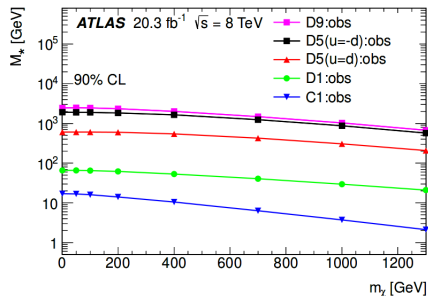
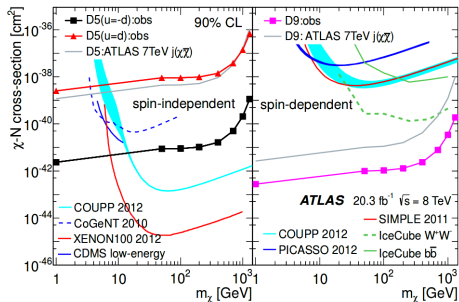
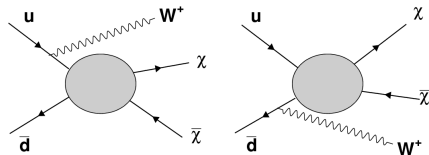


Why MadGraph5?

- Popular among experimentalists.
- Novel Python libraries make add-ons easy.
- MadGraph5_aMC@NLO is an NLO generator (future loop-induced processes).

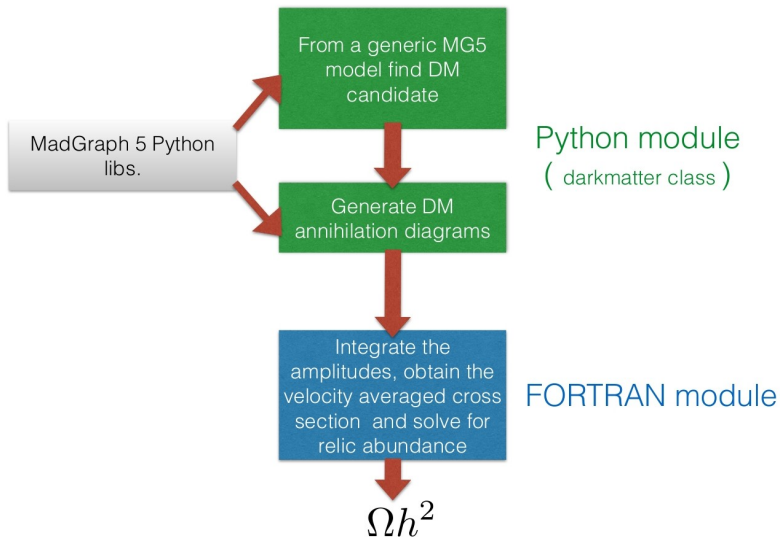
ATLAS search for DM

ATLAS results from a DM search in fat jet + MET channel - constraints on a myriad of effective models.



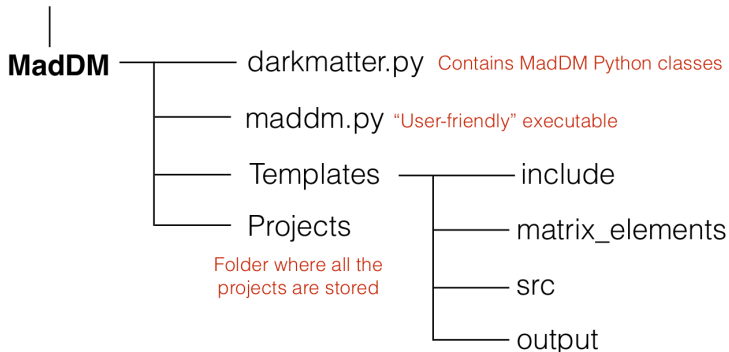
ATLAS used MG for collider analysis, why not add the MadDM output for relic density?

MadDM structure



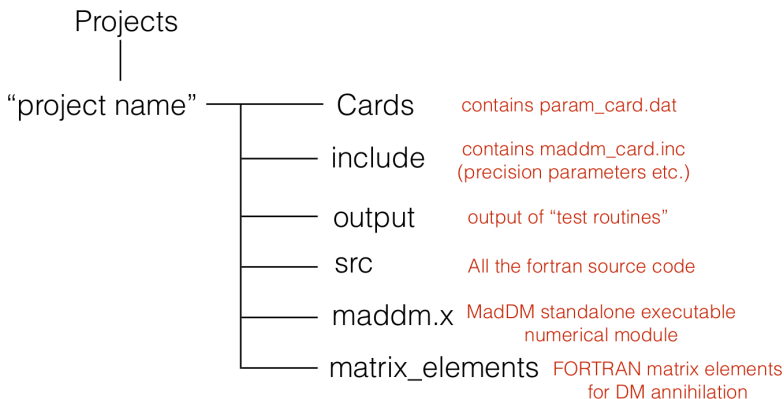
MadDM structure

MadGraph 5



MadDM works out of the box!
 Just unpack the tarball in a
 MadGraph5 folder.

MadDM structure



A project folder is a result of executing the Python MadDM module. It is standalone and does not require any further Python MadDM code

example.py

```
#!/usr/bin/env python
from init import *
from darkmatter_eff import *

#Create the relic density object.
dm=darkmatter()
#Initialize it from the rsxSM model in the MadGraph model folder,
#and store all the results in the Projects/rsxSM subfolder.
dm.init_from_model('rsxSM', 'rsxSM_project', new_proj = True)
#dm.init_from_model('DM_eff_scalar', 'DM_eff_scalar', new_proj = True)

# Determine the dark matter candidate...
dm.FindDMCandidate(prompts=False, dm_candidate='')

#...and all the coannihilation partners with the mass splitting
# defined by  $|m_{X0} - m_{X1}| / m_{X0} < \text{coann\_eps}$ .
dm.FindCoannParticles(prompts = False, coann_eps = 0.1)

#Get the project name with the set of DM particles and see
#if it already exists.
dm.GetProjectName()
```


example.py

```

#Generate all 2-2 diagrams.
print "Generating annihilation diagrams..."
dm.GenerateDiagrams()

#Generate the diagrams for direct detection.
print "Generating direct detection diagrams..."
dm.GenerateDiagramsDirDetect()    (NEW feature !!!)

#Print some dark matter properties in the mean time.
print "----- Testing the darkmatter object properties -----"
print "Calc. Omega: "+str(dm._do_relic_density)
print "Calc. DD: "+str(dm._do_direct_detection)
print "DM name: "+dm._dm_particles[0].get('name')
print "DM spin: "+str(dm._dm_particles[0].get('spin'))
print "DM mass var: "+dm._dm_particles[0].get('mass')
print "Mass: "+ str(dm.GetMass(dm._dm_particles[0].get('pdg_code')))+"\n"
print "Project: "+dm._projectname

#Output the FORTRAN version of the matrix elements
#and compile the numerical code.
dm.CreateNumericalSession()

#Calculate relic density.
omega = dm.CalculateRelicAbundance()
print "-----"
print "Relic Density: "+str(omega),
print "-----"

```

MadDM v1.0 : relic abundance



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MadDM - Relic density calculation

In DM model with only one DM particle, the density evolution is described by the rate equation

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -2 \langle \sigma_{\text{eff}} v \rangle \left(n_\chi^2 - \left(n_\chi^{\text{EQ}} \right)^2 \right),$$

with

$\langle \sigma_{\text{eff}} v \rangle \equiv \sum_{i,j=1}^N \langle \sigma(\chi_i \chi_j \rightarrow SM) v \rangle \frac{n_{\chi_i}^{\text{EQ}} n_{\chi_j}^{\text{EQ}}}{\left(n_\chi^{\text{EQ}} \right)^2}$, the thermally averaged cross section
 n_χ^{EQ} , the equilibrium density

To obtain DM relic abundance in canonical models, integrate the rate equation from a freeze out time (determined by iteration over the rate equation integration)

$$\Omega h^2 \sim \left(\int_{x_f}^{\infty} dx \frac{\langle \sigma v \rangle}{x^2} \right)^{-1} \quad x \equiv \frac{m_\chi}{T}$$

MadDM takes co-annihilation, threshold effects and resonances into account and is also able to deal with non-canonical models !

MadDM - Relic density validation

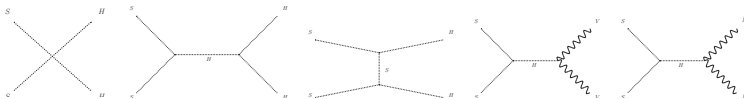
- Toy model : Real singlet extension of SM (rsxSM):

$$\mathcal{L}_{DM} = \frac{m^2}{2} H^\dagger H + \frac{\lambda}{4} (H^\dagger H)^2 + \frac{\delta}{2} H^\dagger H S^2 + \frac{m_S^2}{2} S^2 + \frac{\lambda_S}{4} S^4$$

Only parameters are mass of DM and coupling to the H boson

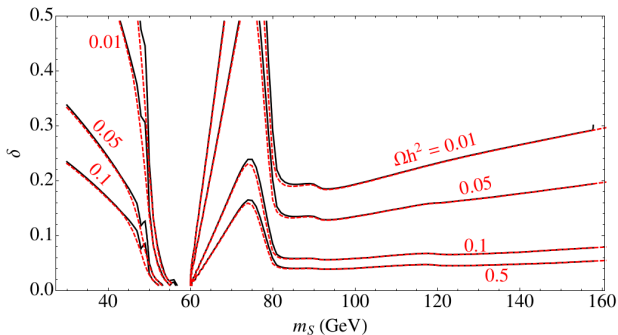


- Typical DM annihilation diagrams :



MadDM - Validations

Real singlet extension of SM



Black - MadDM
Red - micrOMEGAs

Excellent agreement
 over a wide range of DM parameters
 (Some precision issues in the
resonance region where the Higgs
sits right above threshold due to a
saddle point in the integrand)

MadDM v2.0 : direct detection and DDM



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K.C. Kong



G. Mohlabeng



A. Para



J. Yoo

MadDM : Direct detection (NEW)

Typically $v_{DM} \sim 300 \text{ km/s} \ll c$

Transfer momentum to nuclei is very low \rightarrow 0-momentum transfer approximation.

Calculation steps :

- 1 Lagrangian: extract spin-independent and spin-dependent operators.
- 2 DM-quarks amplitude computation with MadGraph.
- 3 DM-nucleon calculation (form factors).
- 4 DM-nucleus interaction.
- 5 Direct detection rates.

MadDM: operator expansion method

We implement an effective Lagrangian (\mathcal{L}_{eff}) including all the effective operators available for DM-quarks interactions for a specific DM spin and for SI (or SD) (similar as micrOMEGAs, arXiv:0803.2360) :

	DM spin	Even	Odd
SI	0	$2M_\chi \Phi_\chi \Phi_\chi^* \bar{\psi}_q \psi_q$	$i \left(\partial_\mu \Phi_\chi \Phi_\chi^* - \Phi_\chi \partial_\mu \Phi_\chi^* \right) \bar{\psi}_q \gamma^\mu \psi_q$
	$\frac{1}{2}$	$\bar{\psi}_\chi \psi_\chi \bar{\psi}_q \psi_q$	$\bar{\psi}_q \gamma_\mu \psi_\chi \bar{\psi}_q \gamma^\mu \psi_q$
	1	$2M_\chi A_{\chi\mu}^* A_{\chi\mu}^H \bar{\psi}_q \psi_q$	$i \lambda_{q,o} \left(A_{\chi\alpha}^{*\alpha} \partial_\mu A_{\chi\alpha} - A_{\chi\alpha}^\alpha \partial_\mu A_{\chi\alpha}^* \right) \bar{\psi}_q \gamma_\mu \psi_q$
SD	$\frac{1}{2}$	$\bar{\psi}_\chi \gamma_\mu \gamma_5 \psi_\chi \bar{\psi}_q \gamma^\mu \gamma_5 \psi_q$	$-\frac{1}{2} \bar{\psi}_\chi \sigma_{\mu\nu} \psi_\chi \bar{\psi}_q \sigma^{\mu\nu} \psi_q$
	1	$\sqrt{6} \left(\partial_\alpha A_{\chi\beta}^* A_{\chi\nu} - A_{\chi\beta}^* \partial_\alpha A_{\chi\nu} \right)$ $\varepsilon^{\alpha\beta\nu\mu} \bar{\psi}_q \gamma_5 \gamma_\mu \psi_q$	$i \frac{\sqrt{3}}{2} \left(A_{\chi\mu} A_{\chi\nu}^* - A_{\chi\mu}^* A_{\chi\nu} \right) \bar{\psi}_q \sigma^{\mu\nu} \psi_q$

Then the trick is to combine \mathcal{L}_{eff} to the input model (related to \mathcal{L}_{input}) to get the interference term between the two models :

$$\left| \mathcal{M}_{eff+input} \right|^2 = \left| \mathcal{M}_{input} \right|^2 + \left| \mathcal{M}_{eff} \right|^2 + 2 \left| \mathcal{M}_{eff} \cdot \mathcal{M}_{input}^* \right|$$

⇒ the interference term gives us the right contribution !

⇒ **Already implemented in MadDM**

DM-nucleus interaction

Matrix element $\langle \bar{q}q \rangle$ of quarks in a nucleon state :

$$\langle \bar{q}q \rangle = \frac{m_{p,n}}{m_q} f_q^{p,n} \text{ (light quarks); } \langle \bar{q}q \rangle = \frac{2}{27} \frac{m_{p,n}}{m_q} f_G^{p,n} \text{ (heavy quarks)}$$

DM-nucleon couplings:

$$f_{p,n}^\chi = m_{p,n} \sum_{q=u,d,s} \frac{C_q}{m_q} f_q^{p,n} + \frac{2}{27} m_{p,n} f_G^{p,n} \sum_{q=c,b,t} \frac{C_q}{m_q}$$

where C_q comes from the projection operator method (and $\langle \mathcal{M} \rangle = C_q \langle \bar{q}q \rangle$).
Finally, you can compute the cross-section DM-nucleus (SI interactions),

$$\sigma = \frac{4m_N^2 m_\chi^2}{\pi (M_\chi + m_N)^2} \cdot (A f_p^\chi + (A - Z) f_n^\chi)^2$$

⇒ Already implemented in MadDM !

DM-nucleus interaction

Spin dependent case :

$$\sigma = \frac{16m_N^2 m_\chi^2}{\pi (M_\chi + m_N)^2} \frac{J_A + 1}{J_A} (\epsilon_p S_p^a + \epsilon_n S_n^A)^2$$

where J_A is the spin of the nucleus, $S_{p,n}^A$ are the expectation value of the spin content of the nucleon in a nucleus with A nucleons

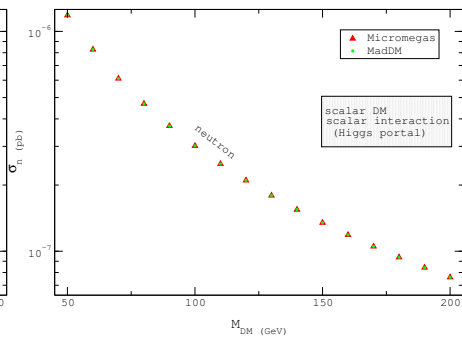
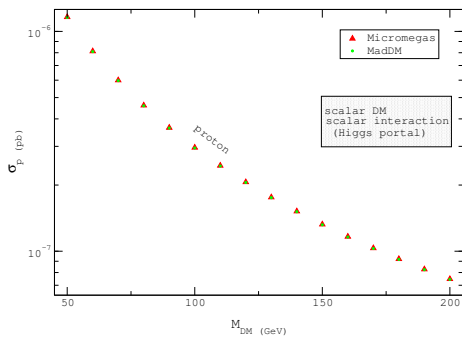
⇒ Already implemented in MadDM !

Differential rate :

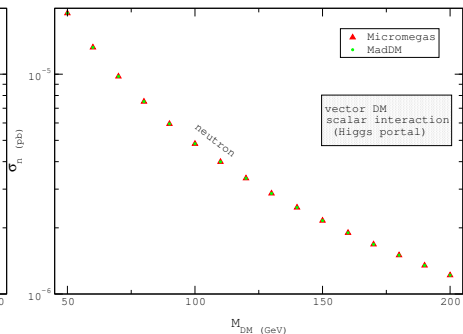
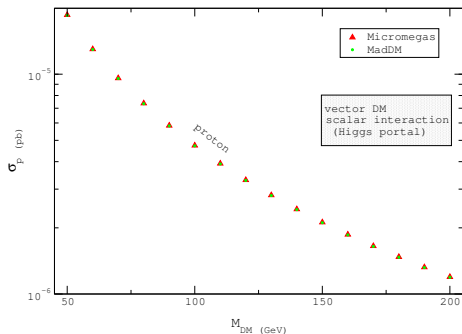
$$\frac{dR}{dE_r} = \frac{\rho_0 \sigma_0}{2m_\chi m_r^2} F(E_r) \int_{v_{min}}^{\infty} \left[\frac{f(v)}{v} \right] dv$$

⇒ Available soon in MadDM !

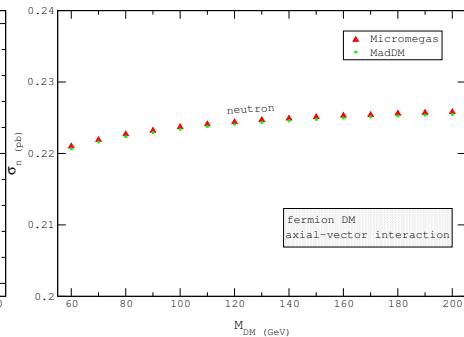
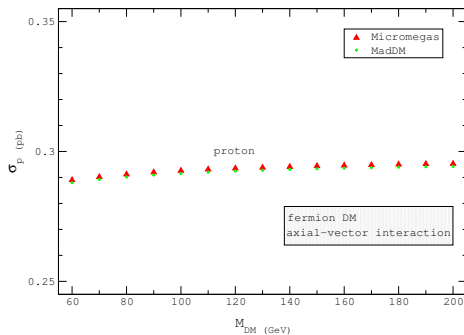
MadDM validations - rsxSM



MadDM validations - Higgs portal (vector DM)



MadDM validations - Axial vector interaction (fermion DM)

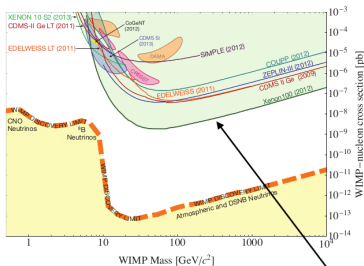


MadDM: directional detection

MadDM - directional detection of DM

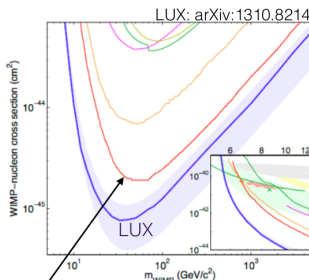
2 big reasons to care about directional information:

- **In case DM is discovered** - need to measure the halo properties. Directional information could be important in this case!
- **In case DM is not discovered** - future 1 ton scale detectors could result in strong limits on WIMP-nucleon cross section. Neutrinos could become a non-negligible background. Directional information can be used to discriminate neutrino backgrounds!

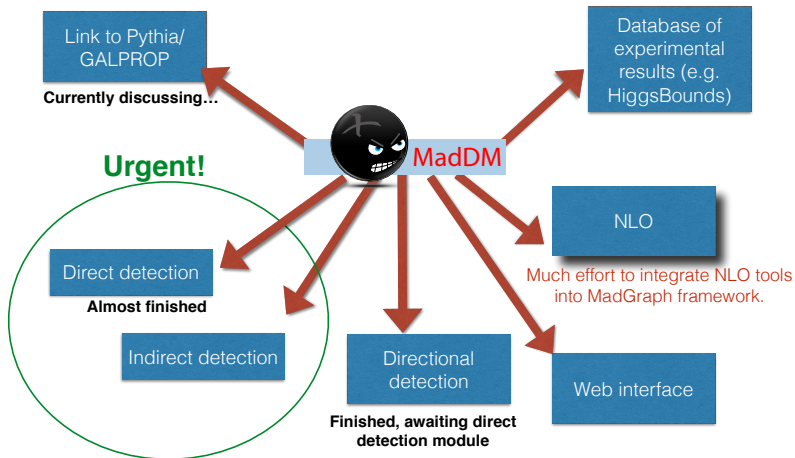


J. Billard, L. Strigari, E. Figueroa-Feliciano
(arXiv:1307.5458)

Xenon100



MadDM - (near) future plans.



susy.phsx.ku.edu/~mihailo/



DEMO !

Tutorial : <http://susy.phsx.ku.edu/~mihailo/tutorial.html>